

The Observer Effect in Estimation with Physical Communication Constraints

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Joint work with

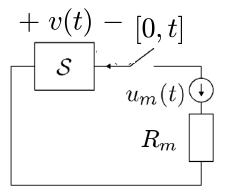
Jean-Charles Delvenne, Univ. catholique de Louvain

IFAC World Congress, Milan Sep. 1, 2011

Observer Effect - Example

Voltage perturbation in passive circuit $S - \Delta v(t)$

Uncertainty in optimal voltage estimate — $\Delta \hat{v}(t)$



With or without active compensation, a trade-off exists:

 $|\Delta v(t)||\Delta \hat{v}(t)| \ge 2k_B T_m/C$ small t

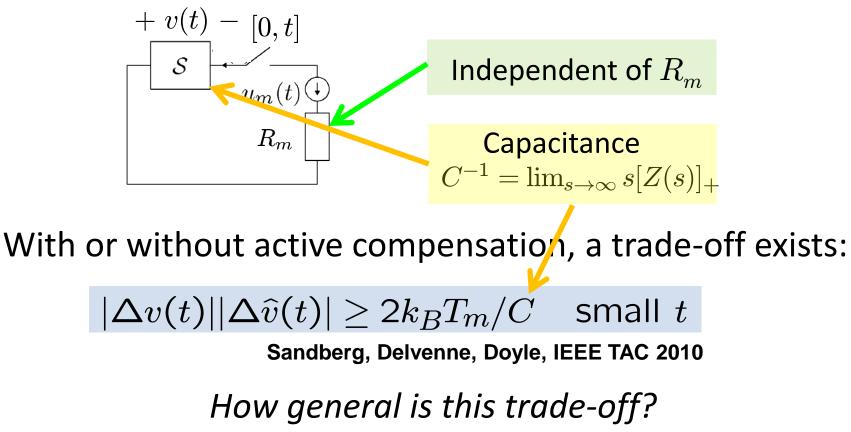
Sandberg, Delvenne, Doyle, IEEE TAC 2010

How general is this trade-off?

Observer Effect - Example

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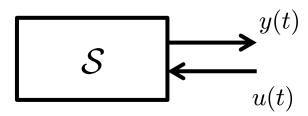


Motivation

- The observer effect has been extensively studied in quantum mechanics, but seems overlooked in a classical mechanics setting
- We want to know the performance limits of devices when resources, such as DOFs, energy, temperature, and time, are finite or limited
- How well and what can we actually implement from a small collection of physical building blocks? Related to circuit synthesis (*M. Smith, B.D.O. Anderson, et al.*)
- Similar questions currently being asked in synthetic biology (*R. Murray, D. Del Vecchio, et al.*)

Measurement Model

 A physical system S with a property y(t) we want to measure. Together with conjugate variable u(t) it forms a port [external work rate = u(t)^Ty(t)]



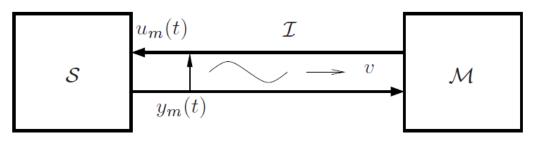
Example: y(t) = velocity, voltage,... u(t) = force, current,...

• Treat u as a *small perturbation* u_m from measurement

$$S \begin{cases} \dot{x}(t) = f(x(t), u_m(t)) \\ y_m(t) = g(x(t)) \\ x(0) = x_0 \end{cases} \approx \begin{cases} \dot{x}(t) = A(t)x(t) + B(t)u_m(t) \\ y_m(t) = C(t)x(t) \\ x(0) = x_0 \end{cases}$$

Measurement Model

Interconnect S to a measurement device M through a communication medium I, during a short time [0,t]



- Problems:
 - What is the best estimate of y(t) we can get?
 - How much must S simultaneously be perturbed?(=back action, retroactivity,...)
- Assumption: Medium *I* is well modeled by a *lossless* wave equation, in thermal equilibrium at time *t*=0

Why Lossless Wave Equation?

Mechanics:

$$u_m(t), y_m(t)$$

- Circuits: $u_m(t), y_m(t)$
- Electromagnetic fields: $\left(\nabla^2 \epsilon \mu \frac{\partial^2}{\partial t^2}\right) E(x,t) = 0$
- At the terminals for times 0 < t < l/v:

$$u_m(t) = -Zy_m(t) - \sqrt{2k_BT_ZZw_Z(t)}$$

(Z>0 [acoustic/characteristic] impedance, T_{z} temperature, v wave velocity, l medium length, and white noise $w_{z}(t)$ due to FD-theorem. Transmission line $Z = \sqrt{\frac{R+i\omega L}{G+i\omega C}} \rightarrow \sqrt{\frac{L}{C}}, \quad \omega \sim 1/t \rightarrow \infty$.)

Where Does the Noise Come From?

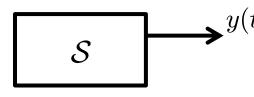
Mechanics:

$$u_m(t), y_m(t)$$

- Circuits: $u_m(t), y_m(t)$
- Electromagnetic fields: $\left(\nabla^2 \epsilon \mu \frac{\partial^2}{\partial t^2}\right) E(x,t) = 0$
- Assume exact initial state of wave equation unknown
- In thermal equilibrium: Assign the expected energy $\frac{1}{2}k_BT_Z$ to each mode
- Total effect at the terminal sums up to white noise, see *Nyquist, 1928* (or *Sandberg, Delvenne, Doyle, 2010*)

Unmeasured vs. Measured System

Unmeasured system • Measured system



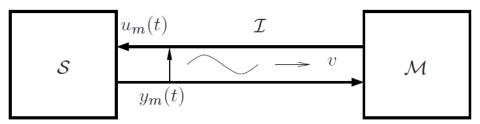
$$\dot{x}(t) = Ax(t) \quad x(0) = x_0,$$

 $y(t) = Cx(t) \quad y_0 := Cx_0$

For small *t*:

$$y(t) = Ce^{At}x_0$$

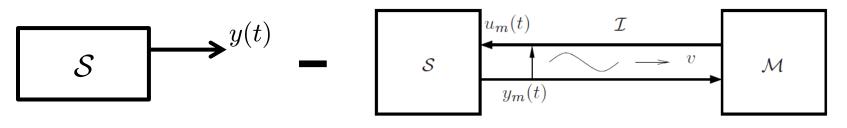
= $y_0 + CAx_0t + O(t^2)$



- $\dot{x}_m(t) = (A BZC)x_m(t) B\sqrt{2k_BT_ZZ}w_Z(t),$ $y_m(t) = Cx_m(t), \quad x_m(0) = x_0,$ $u_m(t) = -Zy_m(t) \sqrt{2k_BT_ZZ}w_Z(t),$
 - For small *t*: $y_m(t) = Ce^{(A - BZC)t} x_0$ $-\int_0^\tau Ce^{(A-BZC)(t-\tau)} B\sqrt{2k_BT_ZZ} w_Z(\tau) d\tau$ $= y_0 + C(A - BZC)x_0t$ $-CB\sqrt{2k_BT_ZZ}\int_0^t w_Z(\tau)d\tau + O(t\sqrt{t}),$

Measurement Back Action

Difference between un-measured and measured output
=: back action of measurement



• Deterministic back action:

$$y(t) - \mathbf{E}y_m(t) = M^{-1}Zy_0t + O(t^2)$$

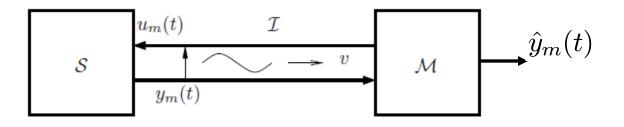
(Present even when $T_z = 0$. $M^{-1} := CB$ is the "inverse inertia" of S)

Stochastic back action:

$$Var[y_m(t)] = 2k_B T_Z M^{-1} Z M^{-T} t + O(t^2)$$

Idealized Measurement Device M*

• The measurement device *M* produces an estimate \hat{y}_m of y_m



• In best case, M has exact models of S and I, and knows the temperature T_z (but not exact state of I) \Rightarrow Kalman filter M^* optimal:

 $\dot{\hat{x}}_{m}(t) = (A - BZC)\hat{x}_{m}(t) + K(t)(u_{m}(t) - \hat{u}_{m}(t))$ $\hat{u}_{m}(t) = -ZC\hat{x}_{m}(t),$ $\hat{y}_{m}^{*}(t) = C\hat{x}_{m}(t),$

• Any M with less or equal knowledge is no better than M^{*}

Lower Bounds on Accuracy

Accuracy of *M*^{*} determined from differential filter Riccati equation:

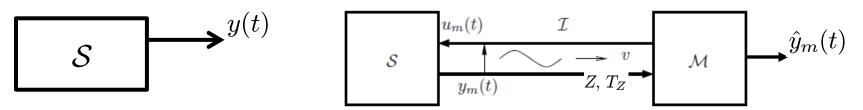
$$\dot{P} = (A - BZC)P + P(A - BZC)^{T}$$

- $(PC^{T}Z - 2k_{B}T_{Z}BZ)(2k_{B}T_{Z}Z)^{-1}(PC^{T}Z - 2k_{B}T_{Z}BZ)^{T}$
+ $2k_{B}T_{Z}BZB^{T}$

- Assume zero knowledge at t = 0: $P(0) = \mathbf{E}[\hat{x}_m(0) x_0]^2 = +\infty$
- Series expansion of P(t) gives $P(t) = P_{-1}t^{-1} + P_0 + P_1t + O(t^2)$ $P_{-1} = P_{-1}C^T Z C P_{-1}/2k_B T_Z$
- Optimal measurement accuracy:

 $\operatorname{Var}[\hat{y}_m^*(t) - y_m(t)] = 2k_B T_Z Z^{-1} t^{-1} + O(1)$

Back Action and Accuracy Trade-Off

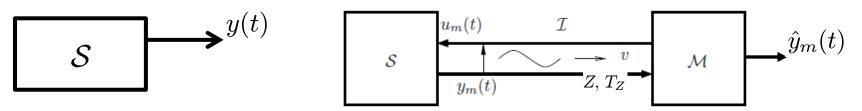


- Back action: $y(t) \mathbf{E}y_m(t) = M^{-1}Zy_0t + O(t^2),$ $\operatorname{Var}[y_m(t)] = 2k_BT_ZM^{-1}ZM^{-T}t + O(t^2)$
- Measurement accuracy: $\operatorname{Var}[\hat{y}_m(t) y_m(t)] \ge \operatorname{Var}[\hat{y}_m^*(t) y_m(t)]$ = $2k_B T_Z Z^{-1} t^{-1} + O(1)$

• Trade-off:
$$\sqrt{\operatorname{Tr} \operatorname{Var}[y_m(t)]} \sqrt{\operatorname{Tr} \operatorname{Var}[\hat{y}_m(t) - y_m(t)]} \ge 2k_B T_Z \operatorname{Tr} M^{-1} + O(t)$$

Trade-off independent on small t and medium impedance Z!

Back Action and Accuracy Trade-Off

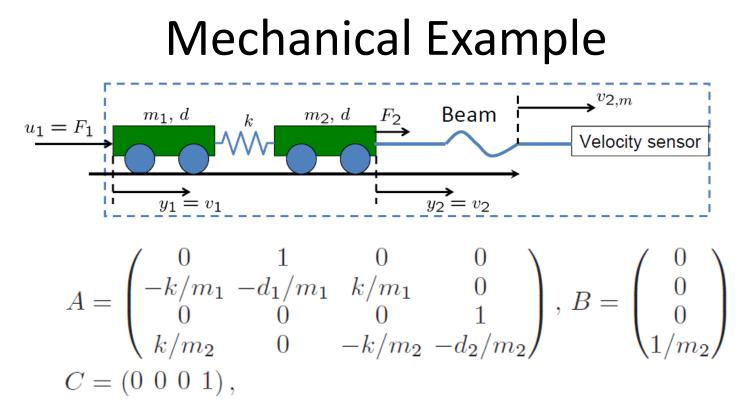


- **Back action:** $y(t) \mathbf{E}y_m(t) = M^{-1}Zy_0t + O(t^2),$ $\operatorname{Var}[y_m(t)] = 2k_BT_ZM^{-1}ZM^{-T}t + O(t^2)$
- Measurement accuracy: $\operatorname{Var}[\hat{y}_m(t) y_m(t)] \ge \operatorname{Var}[\hat{y}_m^*(t) y_m(t)]$ = $2k_B T_Z Z^{-1} t^{-1} + O(1)$

back action \times accuracy

• Trade-off: $\sqrt{\operatorname{Tr} \operatorname{Var}[y_m(t)]} \sqrt{\operatorname{Tr} \operatorname{Var}[\hat{y}_m(t) - y_m(t)]} \ge 2k_B T_Z \operatorname{Tr} M^{-1} + O(t)$

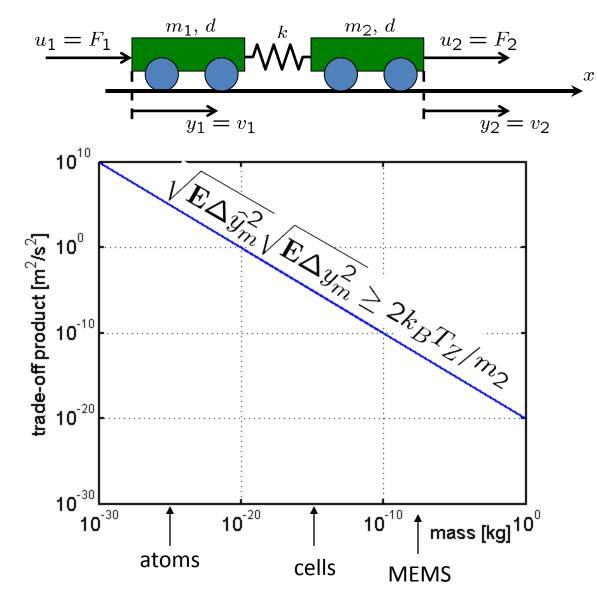
Trade-off independent on small t and medium impedance Z!



- Inverse inertia $M^{-1}:=CB=1/m_2$
- Det. back action $Zv_{2,0}t/m_2$, stoch. back action $2k_BT_ZZt/m_2^2$
- Measurement accuracy $2k_BT_Z/Zt$
- Trade-off: back action × accuracy $\geq 2k_BT_Z/m_2$

Mechanical Example

- $k_B = 1.4 \cdot 10^{-23} \text{ J/K}$
- *T_z* = 300 K
- Not visible at macroscopic level
- Significant at microscopic level



Outline

- Measurement model
- Measuring deterministic systems
- Measuring port-Hamiltonian systems in thermal equilibrium (details in paper)

System S in Thermal Equilibrium

back action
$$\geq \frac{2T_Z}{T_S} \times \Delta \text{accuracy}$$

- Only ratio T_Z/T_S determines trade-off. Is this result essentially different from before? back action × accuracy $\ge 2k_BT_Z/m$
- Not really, as system temperature $T_{\rm S}\to\infty$ we obtain earlier result since rate of learning also goes to infinity
- System inertia disappeared and was replaced by system temperature $T_{\rm S}$
- When $T_{\rm S}$ is really small, only small improvement in accuracy since we already know a lot about the system

Summary

- There are general trade-offs between measurement accuracy and back action for classical measurements
- Holds not only for passive system *S* and particular *M* (compare *Sandberg*, *Delvenne*, *Doyle*, *IEEE TAC 2010*)
- Lossless measurement medium of temperature $T_{\rm Z}$
- System with "inverse inertia" $M^{-1} = CB$ (first Markov parameter) and completely unknown initial state back action × accuracy $\geq 2k_BT_Z \operatorname{Tr} M^{-1}$
- System of temperature $T_{\rm S}$

back action
$$\geq \frac{2T_Z}{T_S} \times \Delta \text{accuracy}$$