## Quantum

## Lecture 12

## - Quantum algorithms

## - Quantum search

## - The quantum Fourier transform

## - Quantum simulation



## Quantum algorithms

$$
\begin{aligned}
& \mathcal{O}\left(g_{n}\right)=\left\{f_{n}: 0 \leq f_{n} \leq c g_{n} \text { for } n \geq n_{0}\right\} \\
& \text { for some } c>0 \text { and integer } n_{0}>0
\end{aligned}
$$

"Complexity $\mathcal{O}\left(g_{n}\right) " \Longleftrightarrow$ true complexity $c_{n} \in \mathcal{O}\left(g_{n}\right)$

## Quantum Search

Generic search problem
For $x \in[0: N-1]$ assume that $f(x)=1$ for $x \in \mathcal{M} \subset[0: N-1]$, $|\mathcal{M}|=M<N(M \ll N)$, and $f(x)=0$ o.w.
$\mathcal{M}$ is the set of solutions to $f(x)$
The problem is to find one solution, i.e. one $x \in \mathcal{M}$
Assume that we have an oracle that can check the value $f(x)$ for one given $x$ at low cost

In general (i.e. not only for search)
$\mathbb{P}=\{$ can be solved with complexity $\mathcal{O}$ (a polynomial) $\}$
$\mathbb{N} \mathbb{P}=\{$ has an oracle of complexity $\mathcal{O}$ (a polynomial) $\}$
Not known if $\mathbb{N P}=\mathbb{P}$

For a basis $\{|x\rangle\}_{x=0}^{N-1}$ the quantum oracle $O$ is the operator

$$
O|x\rangle=(-1)^{f(x)}|x\rangle
$$

The Grover operator
$G|x\rangle=(2|\psi\rangle\langle\psi|-I) O|x\rangle$
Assume $N=2^{n}$ and let

$$
|\psi\rangle=2^{-\frac{n}{2}} \sum_{x=0}^{N-1}|x\rangle
$$

where each $|x\rangle$ corresponds to $n$ qubits $(|0\rangle=|00 \cdots 0\rangle$ etc.)

Let $\mathcal{N}=[0: N-1] \backslash \mathcal{M}$ and

$$
|\alpha\rangle=\frac{1}{\sqrt{N-M}} \sum_{x \in \mathcal{N}}|x\rangle, \quad|\beta\rangle=\frac{1}{\sqrt{M}} \sum_{x \in \mathcal{M}}|x\rangle
$$

If we define

$$
\cos \frac{\theta}{2}=\sqrt{\frac{N-M}{N}} \Rightarrow \sin \frac{\theta}{2}=\sqrt{\frac{M}{N}}
$$

then

$$
|\psi\rangle=\cos \frac{\theta}{2}|\alpha\rangle+\sin \frac{\theta}{2}|\beta\rangle
$$

and

$$
G^{k}|\psi\rangle=\cos \left(\frac{2 k+1}{2} \theta\right)|\alpha\rangle+\sin \left(\frac{2 k+1}{2} \theta\right)
$$



Each time $G$ is applied, the initial state $|\psi\rangle$ is taken closer to $|\beta\rangle$
Quantum search (for $M<N / 2$ ): Prepare the state $|\psi\rangle$
Iterate the Grover operator $K$ times
Measure $\Rightarrow$ a state $|x\rangle^{\prime} \in\{|x\rangle: x \in \mathcal{M}\}$ with high probability For $M \ll N$ choosing $K=\lceil\sqrt{N / M}\rceil$ gives a probability of success of at least $1-M / N$

## The Quantum Fourier Transform

Assume $\mathcal{H}$ is $N$-dimensional, and let $\{|k\rangle\}_{k=0}^{N-1}$ be a basis. For an arbitrary state $|\psi\rangle=\sum_{k} x_{k}|k\rangle$, let $\mathcal{F}$ be the operator defined by

$$
\mathcal{F}|\psi\rangle=\sum_{k} y_{k}|k\rangle
$$

where

$$
y_{k}=\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} x_{j} e^{2 \pi i j k / N}
$$

is the discrete Fourier transform of $\left\{x_{j}\right\}$
$\mathcal{F}|\psi\rangle$ is the quantum Fourier transform of $|\psi\rangle$
$\mathcal{F}$ is a unitary transformation

Assume that $N=2^{n}$ for some integer $n$, and for $j \in[0: N-1]$ let

$$
j=\sum_{\ell=1}^{n} j_{\ell} 2^{n-\ell}
$$

be the binary expansion of $j$ in terms of $\left\{j_{\ell}\right\}, j_{\ell} \in\{0,1\}$
Define the notation

$$
j=j_{1} j_{2} \cdots j_{n}=\sum_{\ell=1}^{n} j_{\ell} 2^{n-\ell} \in[0: N-1]
$$

and, for $1 \leq k \leq \ell \leq n$,

$$
0 . j_{k} j_{k+1} \cdots j_{\ell}=\sum_{i=k}^{\ell} j_{i} 2^{k-i-1} \in[0,1)
$$

Identify $\{|j\rangle\}$ with an $n$-fold qubit basis via $|j\rangle \leftrightarrow\left|j_{1} \cdots j_{n}\right\rangle$
Then we can write $\mathcal{F}\left|j_{1} \cdots j_{n}\right\rangle=$
$2^{-\frac{n}{2}}\left(|0\rangle+e^{2 \pi i 0 . j_{n}}|1\rangle\right)\left(|0\rangle+e^{2 \pi i 0 . j_{n-1} j_{n}}|1\rangle\right) \cdots\left(|0\rangle+e^{2 \pi i 0 \cdot j_{1} \cdots j_{n-1} j_{n}}|1\rangle\right)$

## Phase estimation

Assume we wish to estimate the eigenvalue $\lambda=e^{2 \pi i \phi}$ corresponding to the eigenvector $|u\rangle$ of a unitary operator $U$

Assume $\phi$ has an exact $t$-bits expansion, $\phi=0 . f_{1} \cdots f_{t}$
If we, without knowing $\phi$, can compute the state
$2^{-\frac{t}{2}}\left(|0\rangle+e^{2 \pi i 0 . f_{t}}|1\rangle\right)\left(|0\rangle+e^{2 \pi i 0 . f_{t-1} f_{t}}|1\rangle\right) \cdots\left(|0\rangle+e^{2 \pi i 0 . f_{1} \cdots f_{t-1} f_{t}}|1\rangle\right)$
then an inverse Fourier transform will result in $\left|f_{1} f_{2} \cdots f_{t}\right\rangle$
A measurement in the qubit basis then gives $\phi$
If $\phi$ is not on the form $0 . f_{1} \cdots f_{t}$ for some $t$, then using

$$
t=n+\left\lceil\log \left(2+\frac{1}{2 \varepsilon}\right)\right\rceil
$$

qubits will give $n$ bits accuracy and error probability $\leq \varepsilon$


Phase estimation: Need to prepare the state $|u\rangle$; Need to implement the $U^{j}$ mappings; Complexity $\mathcal{O}\left(t^{2}\right)$

## Order finding

Greatest common divisor of a set $A$ of integers = biggest integer that divides all numbers in the set, notation $\operatorname{gcd}(A)$
Two integers $q_{1}$ and $q_{2}$ are relatively prime (coprime) if $\operatorname{gcd}\left(q_{1}, q_{2}\right)=1$

The order $r$ of an integer $x$ modulo a prime number $p$ is the smallest integer $r$ such that $x^{r}=1 \bmod p$
Finding $r$ is believed to be hard on a classical computer, in the sense that the complexity is at least linear in $p$,

Fermat's little theorem: $x^{p-1}=1 \bmod p \Rightarrow r<p$
Order of $x$ modulo a non-prime $M: x^{\varphi(M)}=1 \bmod M$ where

$$
\varphi(M)=|\{y: 1 \leq y \leq M, \operatorname{gcd}(y, M)=1\}|
$$

i.e., the complexity is still linear in $M$

Defining the unitary operation $U$ as $U|y\rangle=|x y \bmod M\rangle$, we have with

$$
\left|u_{s}\right\rangle=\frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2 \pi i s k / r}\left|x^{k} \bmod M\right\rangle
$$

for $0 \leq s \leq r-1$, that

$$
U\left|u_{s}\right\rangle=e^{2 \pi i s / r}\left|u_{s}\right\rangle
$$

Phase estimation $\Rightarrow\left\{e^{2 \pi i s / r}\right\} \Rightarrow r$ with complexity $\mathcal{O}\left((\log M)^{3}\right)$
We need $r$ to prepare $\left|u_{s}\right\rangle$ : Can use $|1\rangle$ instead of $\left|u_{s}\right\rangle$, since

$$
\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1}\left|u_{s}\right\rangle=|1\rangle
$$

## Factoring

Prime factoring: Given a (large) positive integer $q$, find a prime number $p$ that divides $q$

Believed to be hard on a classical computer, with complexity $\mathcal{O}(\sqrt{q})$ - The factoring problem being "hard" is a crucial assumption in public key encryption

Assume $q$ is odd (otherwise 2 is a trivial factor)
For $x \in[2: q-2]$ suppose $x^{2}=1 \bmod q$. Then at least one of $\operatorname{gcd}(x-1, q)$ and $\operatorname{gcd}(x+1, q)$ is a factor in $q$
Suppose $q$ has $m$ different prime factors and let $x$ be an integer chosen uniformly in $[1: q-1] \cap\{s: s$ and $q$ relatively prime $\}$, then

$$
\operatorname{Pr}\left(r \text { is even and } x^{\frac{r}{2}} \neq-1 \bmod q\right) \geq 1-\frac{1}{2^{m}}
$$

where $r$ is the order of $x \bmod q$

Algorithm: Given an odd number $q>1$
Check if $q=a^{b}$ for some prime $a$ and integer $b$
Choose $x$ at random in $[1: q-1]$; if $\operatorname{gcd}(x, q)>1$ return $\operatorname{gcd}(x, q)$
Use quantum order finding to find the order $r$ of $x \bmod q$
If $r$ is even and $x^{r / 2} \neq-1 \bmod q$ then compute $\operatorname{gcd}\left(x^{r / 2}-1, q\right)$ and $\operatorname{gcd}\left(x^{r / 2}+1, q\right)$ and check if one of these is a factor Otherwise terminate with an error

The steps performed using classical computing have complexity $\mathcal{O}\left((\log q)^{3}\right)$, so the overall complexity relies on the order finding

## Quantum Simulation

Classical system with state in $\mathbb{R}^{d}$ : In general, complexity of simulation grows as $\mathcal{O}(d)$
$N$ quantum particles with states in $\mathcal{H}$ of dimension $d$, complexity of simulating the combined system is in general $\mathcal{O}\left(d^{N}\right)$

Assume $N$ interacting sub-systems such that the evolution of the joint system is described by

$$
i \frac{d}{d t}|\psi\rangle=H|\psi\rangle \Rightarrow|\psi(t)\rangle=e^{-i H t}|\psi(0)\rangle
$$

with $H$ of the form

$$
H=\sum_{\ell=1}^{L} H_{\ell}
$$

where $L=\mathcal{O}(N)$ and each $H_{\ell}$ acts only on few subsystems

Assume the action of each $H_{\ell}, \exp \left(-i H_{\ell} t\right)$, can be simulated efficiently on a quantum computer
We get

$$
|\psi(t)\rangle=e^{-i H t}|\psi(0)\rangle
$$

where we can use the Trotter formula

$$
\lim _{n \rightarrow \infty}\left(e^{\frac{i A t}{n}} e^{\frac{i B t}{n}}\right)^{n}=e^{i(A+B) t}
$$

(for $A$ and $B$ self-adjoint/Hermitian)

## Quantum simulation:

For subsystems of dimension $\mathcal{O}(d)$, the total dimension is $\mathcal{O}\left(d^{N}\right)$
Approximate each $H_{\ell}$ at resolution $\mathcal{O}\left(N^{k}\right)$ (some $k \geq 1$ ) qubits
Simulate each subsystem on a quantum computer
Combine using Trotter's formula, or similar

