# Quantum

Lecture 7

- The Holevo bound
- Typical sequences and subspaces
- Compression

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The Holevo Bound

Assume a discrete random variable  $X \in \mathcal{X}$  with pmf p(x) is embedded on a set of states  $\rho_x$ , as the ensemble  $\{p(x), \rho_x\}$ 

A measurement described by  $\{M_n\}_{n=1}^N$  is performed, resulting in  $Y \in \{1,\ldots,N\}$ 

The Holevo bound states that

$$I(X;Y) \le S(\rho) - \sum_{x \in \mathcal{X}} p(x)S(\rho_x)$$

over all possible  $\{M_n\}$ , and with

$$\rho = \sum_{x \in \mathcal{X}} p(x) \rho_x$$

The entity

$$\chi(p(x), \rho_x) = S(\rho) - \sum_{x \in \mathcal{X}} p(x)S(\rho_x)$$

is the Holevo information of the ensemble  $\{p(x), \rho_x\}$ Note that the joint entropy of the classical-quantum state

$$\sigma = \sum_{x \in \mathcal{X}} p(x) |e(x)\rangle \langle e(x)| \otimes \rho_x$$

(where  $\{e(x)\}$  is a basis) is  $H(p) + \sum_{x \in \mathcal{X}} p(x)S(\rho_x)$ , hence

$$\chi(p(x), \rho_x) = H(p) + S(\rho) - S(\sigma)$$

= mutual information between the classical and the quantum state

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## Fano's Inequality

For discrete random variables, consider

X =variable of interest

Y = observed variable

 $\hat{X} = f(Y)$  estimate of X based on Y

With  $P_e = \Pr(\hat{X} \neq X)$  and  $h(x) = -x \log x - (1-x) \log(1-x)$ , we have Fano's inequality

$$h(P_e) + P_e \log(|\mathcal{X}| - 1) \ge H(X|Y)$$

Hence, in the quantum setting:

For any measurement that tries to conclude X as  $\hat{X}$  from  $\rho$ ,

 $h(P_e) + P_e \log(|\mathcal{X}| - 1) \ge H(X) - S(\rho) + \sum_{x \in \mathcal{X}} p(x)S(\rho_x)$ 

### **Typical Sequences**

For a sequence  $x^n = (x_1, \ldots, x_n)$  with letters in  $\mathcal{X}$  and a pmf p(x) on  $\mathcal{X}$ , let

$$T(x^n) = -\frac{1}{n} \sum_{i} \log p(x_i)$$

For fixed n and  $\varepsilon > 0$ , let

$$\mathcal{T}_{\varepsilon}^{(n)} = \{x^n : |T(x^n) - H(p)| \le \varepsilon\}$$

be the set of  $\varepsilon$ -typical sequences (of length n, given p)

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By the (weak) LLN, if  $X^n \sim \prod_i p(x_i)$  then for any  $\varepsilon > 0$  there is an N such that for all n > N

$$\Pr(X^n \in \mathcal{T}_{\varepsilon}^{(n)}) > 1 - \varepsilon$$

We also have

$$|\mathcal{T}_{\varepsilon}^{(n)}| \le 2^{n(H(p)+\varepsilon)}$$

and there is an N such that for  $n \geq N$ 

$$|\mathcal{T}_{\varepsilon}^{(n)}| \ge (1-\varepsilon)2^{n(H(p)-\varepsilon)}$$

#### Compression

We can enumerate all elements of  $\mathcal{T}_{\varepsilon}^{(n)}$  using numbers from  $[1:M_n]$  with  $M_n \geq \lceil 2^{n(H(p)+\varepsilon)} \rceil$ 

Assume  $X^n \sim \prod_i p(x_i)$ 

Compression code: Observe  $X^n = x^n$ ; if  $x^n \in \mathcal{T}_{\varepsilon}^{(n)}$  then produce  $i \in [1:M_n]$  corresponding to  $x^n$ ; if  $x^n \notin \mathcal{T}_{\varepsilon}^{(n)}$  then declare error

For any  $\varepsilon > 0$ , there is an N such that for all n > N,  $\Pr(\text{error}) \leq \varepsilon$  as long as

$$\frac{1}{n}\log M_n \ge H(p) + \varepsilon + \frac{1}{n}$$

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On the other hand, from Fano's inequality

$$\Pr(\text{error})\frac{\log M_n}{n} + \frac{1}{n} \ge H(p) - \frac{1}{n}\log M_n$$

Hence, for large n, choosing  $n^{-1}\log M_n$  slightly bigger than H(p) is the best compression we can accomplish

#### Preservation of Entanglement

For discrete random variables X and Y with join pmf p(x, y), the mutual information I(X;Y) measures the degree of mutual dependence, or (nonlinear) correlation

In quantum systems, two states are dependent on each-other if they are entangled

Consider a mixed state  $\rho$  in  $\mathcal{H}$  with purification  $|\psi\rangle$  in  $\mathcal{H} \otimes \mathcal{R}$ , i.e.  $\rho = \text{Tr}_{\mathcal{R}} |\psi\rangle \langle \psi|$  for some space  $\mathcal{R}$ 

 $\mathcal{R}$  can model the unknown environment; if we had access to both  $\mathcal{H}$  and  $\mathcal{R}$  then we would be considering the pure state  $|\psi\rangle\langle\psi|$ 

The system  $\mathcal{H}$  is entangled with the environment  $\mathcal{R}$ , as characterized by the entangled state  $|\psi\rangle \in \mathcal{H} \otimes \mathcal{R}$ 

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Assume  $\mathcal{E}$  is applied to  $\rho$  in  $\mathcal{H}$ , resulting in the state  $\sigma$  in  $\mathcal{H} \otimes \mathcal{R}$ . Then, the entanglement fidelity of  $(\rho, \mathcal{E})$  is defined as

$$F(\rho, \mathcal{E}) = \langle \psi | \sigma | \psi \rangle$$

 $F(\rho, \mathcal{E})$  does not depend on  $\mathcal{R}$ ,  $0 \leq F(\rho, \mathcal{E}) \leq 1$ We can easily verify that

$$F(\rho, \mathcal{E}) = (F(|\psi\rangle\langle\psi|, \sigma))^2$$

where  $F(|\psi\rangle\langle\psi|,\sigma)$  is the regular (static) fidelity between the pure state  $|\psi\rangle\langle\psi|$  and  $\sigma$  (remember  $F(\rho,\sigma) = \text{Tr}\sqrt{\rho^{1/2}\sigma\rho^{1/2}}$ )  $F(\rho,\mathcal{E})$  measures how well entanglement is preserved by  $\mathcal{E}$ Let  $\{E_i\}$  be the operation elements of  $\mathcal{E}$ , then we also have

$$F(\rho, \mathcal{E}) = \sum_{i} |\operatorname{Tr}(\rho E_i)|^2$$

#### **Typical Subspaces**

Any density operator  $\rho$  associated with a system  $\mathcal{H}$  has an eigen-decomposition  $\rho = \sum_i \lambda_i |x_i\rangle \langle x_i|$ 

Since  $\sum_i \lambda_i = 1$ , we can interpret this representation for  $\rho$  as an information source;  $|x_i\rangle$  is emitted with probability  $p(x_i) = \lambda_i$ 

Let  $\rho^n = \rho \otimes \cdots \otimes \rho$ ,  $|x^n\rangle = |x_{i_1} \cdots x_{i_n}\rangle = |x_{i_1}\rangle \otimes \cdots \otimes |x_{i_n}\rangle$  and  $\mathcal{H}^n = \mathcal{H} \otimes \cdots \otimes \mathcal{H}$  (*n* times)

The states  $\rho^n$  and  $|x^n\rangle$  correspond to "using the information source"  $(\rho, \mathcal{H})$  a number of n independent times

With 
$$T(|x^n\rangle) = -n^{-1} \sum_{m=1}^n \log p(x_{i_m})$$
 let  
$$\mathcal{T}_{\varepsilon}^{(n)} = \{|x^n\rangle : |T(|x^n\rangle) - S(\rho)| \le \varepsilon$$

and define the typical subspace

$$\mathcal{S}^{(n)}_arepsilon = {
m span}\,\mathcal{T}^{(n)}_arepsilon = {
m span}\{|x^n
angle: |x^n
angle\in\mathcal{T}^{(n)}_arepsilon\}$$

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Let  $P_{\varepsilon}^{(n)}$  denote the projection operator from  $\mathcal{H}^n$  to  $\mathcal{S}_{\varepsilon}^{(n)}$ For any  $\varepsilon > 0$  there is an N such that for n > N

$$\operatorname{Tr}(P_{\varepsilon}^{(n)}\rho^n) \ge 1 - \varepsilon$$

Furthermore, for any n and  $\varepsilon$ 

$$\operatorname{Tr} P_{\varepsilon}^{(n)} \leq 2^{n(S(\rho) + \varepsilon))}$$

and for any  $\varepsilon > 0$  there is an N such that for n > N

$$\operatorname{Tr} P_{\varepsilon}^{(n)} \ge (1-\varepsilon)2^{n(S(\rho)-\varepsilon))}$$

#### Compression



 $\mathcal{C}^n$  maps states in  $\mathcal{H}^n$  to states in a space  $\mathcal{G}_n$  of dimension  $D_n$  $\mathcal{D}^n$  maps states in  $\mathcal{G}_n$  back to states in  $\mathcal{H}^n$ Assume  $|\psi\rangle$  is a purification of  $\rho^n$  in  $\mathcal{H}^n \otimes \mathcal{R}$ , and let  $\mathcal{E}^n = \mathcal{D}^n \circ \mathcal{C}^n$ Let  $\sigma^n$  be the resulting state in  $\mathcal{H}^n \otimes \mathcal{R}$ The corresponding entanglement fidelity is

$$F(\rho^n, \mathcal{E}^n) = \langle \psi | \sigma^n | \psi \rangle$$

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A compression scheme

Select  $\mathcal{G}_n \supset \mathcal{S}_{\varepsilon}^{(n)} \Rightarrow \operatorname{Tr} P_{\varepsilon}^{(n)} \leq D_n$ Set  $\mathcal{C}^n = P_{\varepsilon}^{(n)}$  and  $\mathcal{D}^n = I$  (identity) Then for any  $\varepsilon > 0$  there is an N such that for n > N

$$F(\rho^n, \mathcal{E}^n) \ge |\operatorname{Tr}(\rho^n P_{\varepsilon}^{(n)})|^2 \ge |1 - \varepsilon|^2 \ge 1 - 2\varepsilon$$

It also holds that

$$\operatorname{Tr} P_{\varepsilon}^{(n)} \le 2^{n(S(\rho) + \varepsilon))}$$

Thus  $F(\rho^n, \mathcal{E}^n) > 1 - 2\varepsilon$  as long as

$$\frac{1}{n}\log D_n > S(\rho) + \varepsilon$$

Converse: It can be shown that, if

$$\lim_{n \to \infty} \frac{1}{n} \log D_n < S(\rho)$$

then  $F(
ho^n, \mathcal{S}^n) 
ightarrow 0$  for any projector  $\mathcal{S}^n$ 

If  $\mathcal{H}$  is *d*-dimensional,  $\mathcal{H}^n$  is  $d^n$ -dimensional; i.e. it takes  $n \log d$  qubits to describe a state in  $\mathcal{H}^n$ 

Then the best compression we can have is from  $\log d$  qubits to  $S(\rho)$  ( $\leq \log d$ ) qubits, per use of the source  $(\rho, \mathcal{H})$ ,

with preserved entanglement  $F(\rho^n, \mathcal{S}^n) \to 1$ 

Since  $1 - \sqrt{F(\rho, \sigma)} \leq V(\rho, \sigma) \leq \sqrt{1 - (F(\rho, \sigma))^2}$  we could also use  $V(\rho, \sigma) = 2^{-1} \text{Tr} |\rho - \sigma|$  as fidelity metric,

 $F(\rho,\sigma) \to 1 \iff V(\rho,\sigma) \to 0$ 

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